# THE GROUNDWATER WAVES PROBLEM IN THE DUPUIT-FORCHHEIMER REGIME

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$$\begin{split} \Theta(x,z) &: \text{ porosity} = \frac{\text{void-space volume}}{\text{total volume}} \leq 1\\ \kappa(x,z) &: \text{ permeability (or hydraulic conductivity)}\\ B(x) &: \text{ bathymetry} \end{split}$$



S(x) : substratum

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▶ Water Resources: baseflow, flooding, salt intrusion, pollutant...



▶ Geomorphology: sediment transport, soil alteration...







▶ Energy dissipation:  $\partial_t \mathscr{P}(\eta) + \nabla \cdot \mathscr{G}(\eta, u, w) = -\widetilde{\mathscr{K}}(\eta, u, w)$ 

where 
$$\mathscr{P}(\eta) := g \int_{S}^{\eta} \Theta z \, dz$$
  
 $\mathscr{G}(\eta, u, w) := \int_{S}^{\eta} (gz + p) \Theta u \, dz$   
and  $\widetilde{\mathscr{K}}(\eta, u, w) := \int_{S}^{\eta} \frac{\Theta}{\kappa} (|u|^{2} + w^{2}) \, dz.$ 

(potential energy)

(energy flux)

(dissipation  $\approx$  kinetic energy)

► The **linearized** groundwater waves model with constant parameters has solutions of the form  $\eta(t,x) = D\cos(k \cdot x)e^{-\frac{g\kappa}{D}\tilde{\lambda}t}$ with the **linear decay rate** reads  $\tilde{\lambda} = |kD| \tanh(|kD|)$ .



▶ Integrating the incompressibility constraint, we get  $\partial_t \mathcal{V}(\eta) + \nabla \cdot (\mathcal{V}(\eta)\overline{u}) = 0$ with the water volume  $\mathcal{V}(\eta) := \int_S^{\eta} \Theta dz$ , and the mean horizontal velocity  $\overline{u} := \frac{1}{\mathcal{V}(\eta)} \int_S^{\eta} \Theta u dz$ .

▶ Integrating the horizontal Darcy law, we get  $\mathcal{V}(\eta)\overline{u} = -\int_{S}^{\eta} \kappa \nabla p \, dz$ 

DUPUIT-FORCHHEIMER ASSUMPTION

$$p(t,x,z) = g(\eta(t,x)-z) + O(\mu) \quad \Rightarrow \quad \overline{u} = -gK(\eta)\nabla\eta + O(\mu) \quad \text{with } K(\eta) = \frac{1}{V(\eta)} \int_{S}^{\eta} \Theta \kappa \, \mathrm{d}z.$$

[Dupuit'1863]

 $\mathbb{Z}$  DUPUIT-FORCHHEIMER MODEL (*DF*):  $O(\mu)$ -approximate model

$$\partial_t \mathcal{V}(\eta) - \nabla \cdot (gK(\eta)\mathcal{V}(\eta)\nabla\eta) = 0$$

► Energy dissipation:  $\partial_t \mathscr{P}(\eta) + \nabla \cdot \mathscr{G}(\eta, u, w) = -\widetilde{\mathscr{K}}(\eta, u, w)$ where  $\mathscr{P}(\eta) := g \int_{S}^{\eta} \Theta z \, dz$ ,  $\mathscr{G}(\eta, u, w) := g^2 \eta K \mathcal{V} \nabla \eta$  and  $\widetilde{\mathscr{K}}(\eta, u, w) := g^2 K \mathcal{V} |\nabla \eta|^2$ . ► Linear decay rate:  $\widetilde{\lambda} = |kD|^2$ .



► Assuming constant coefficients:  $\Theta(x,z) = \overline{\Theta}(x) + O(\mu^2)$  and  $\kappa(x,z) = \overline{\kappa}(x) + O(\mu^2)$ ► Assuming the hydrodynamic pressure small:  $p(t,x,z) - g(\eta(t,x) - z) = O(\mu) = \mu q(t,x,z)$ 

- **Assuming** constant coefficients:  $\Theta(x,z) = \overline{\Theta}(x) + O(\mu^2)$  and  $\kappa(x,z) = \overline{\kappa}(x) + O(\mu^2)$  **Assuming** the hydrodynamic pressure small:  $p(t,x,z) g(\eta(t,x) z) = O(\mu) = \mu q(t,x,z)$ ►
- Form the horizontal Darcy law

Form the horizontal Darcy law:  

$$\begin{aligned} u &= -\overline{\kappa}\nabla(g\eta + \mu q) \\ &\text{we get } h\overline{u} = -g\overline{\kappa}h\nabla(h+S) - \mu\overline{\kappa}(\nabla(h\overline{q}) + \beta q_{S}\nabla S) + O(\mu^{2}) \end{aligned}$$
with  $\overline{q}(t,x) = \frac{1}{h(t,x)}\int_{S}^{\eta(t,x)}q(t,x,z)\,\mathrm{d}z$  and  $q_{S}(t,x) = q(t,x,S(x)).$ 

- ► Assuming constant coefficients:  $\Theta(x,z) = \overline{\Theta}(x) + O(\mu^2)$  and  $\kappa(x,z) = \overline{\kappa}(x) + O(\mu^2)$ ► Assuming the hydrodynamic pressure small:  $p(t,x,z) - g(\eta(t,x) - z) = O(\mu) = \mu q(t,x,z)$
- ► Form the horizontal Darcy law:  $u = -\overline{\kappa}\nabla(g\eta + \mu q)$ we get  $h\overline{u} = -g\overline{\kappa}h\nabla(h+S) - \mu\overline{\kappa}(\nabla(h\overline{q}) + \beta q_S\nabla S) + O(\mu^2)$ with  $\overline{q}(t,x) = \frac{1}{h(t,x)}\int_{S}^{\eta(t,x)}q(t,x,z) dz$  and  $q_S(t,x) = q(t,x,S(x))$ .
- ▶ Form the incompressibility constraint:

we get 
$$w(t,x,z) = \overline{w}(t,x) + \frac{z-S-\frac{h}{2}}{h} \widetilde{w}(t,x) + O(\mu^2)$$
 with  $\overline{w} = \mu \left(\beta \overline{u} \cdot \nabla S - \frac{h}{2\overline{\Theta}} \nabla \cdot \left(\overline{\Theta} \overline{u}\right)\right)$   
and  $\widetilde{w} = -\frac{h}{\overline{\Theta}} \nabla \cdot \left(\overline{\Theta} \overline{u}\right)$ 

 $\nabla \cdot \left(\overline{\Theta}u\right) + \partial_z \left(\overline{\Theta}w\right) = 0$ 

- ► Assuming constant coefficients:  $\Theta(x,z) = \overline{\Theta}(x) + O(\mu^2)$  and  $\kappa(x,z) = \overline{\kappa}(x) + O(\mu^2)$ ► Assuming the hydrodynamic pressure small:  $p(t,x,z) - g(\eta(t,x) - z) = O(\mu) = \mu q(t,x,z)$
- ► Form the horizontal Darcy law:  $u = -\overline{\kappa}\nabla(g\eta + \mu q)$ we get  $h\overline{u} = -g\overline{\kappa}h\nabla(h+S) - \mu\overline{\kappa}(\nabla(h\overline{q}) + \beta q_S\nabla S) + O(\mu^2)$ with  $\overline{q}(t,x) = \frac{1}{h(t,x)}\int_{S}^{\eta(t,x)}q(t,x,z) dz$  and  $q_S(t,x) = q(t,x,S(x))$ .
- ► Form the incompressibility constraint: we get  $w(t,x,z) = \overline{w}(t,x) + \frac{z - S - \frac{h}{2}}{h} \widetilde{w}(t,x) + O(\mu^2)$ with  $\overline{w} = \mu \left(\beta \overline{u} \cdot \nabla S - \frac{h}{2\overline{\Theta}} \nabla \cdot \left(\overline{\Theta} \overline{u}\right)\right)$ and  $\widetilde{w} = -\frac{h}{\overline{\Theta}} \nabla \cdot \left(\overline{\Theta} \overline{u}\right)$
- ▶ Form the vertical Darcy law:

$$\begin{split} w &= -\overline{\kappa}\partial_z q\\ \text{we get } h\overline{w} &= \overline{\kappa}q_S \quad \text{and} \quad h\widetilde{w} = 6\overline{\kappa}\left(2\overline{q} - q_S\right) \end{split}$$

- ► Assuming constant coefficients:  $\Theta(x,z) = \overline{\Theta}(x) + O(\mu^2)$  and  $\kappa(x,z) = \overline{\kappa}(x) + O(\mu^2)$ ► Assuming the hydrodynamic pressure small:  $p(t,x,z) - g(\eta(t,x) - z) = O(\mu) = \mu q(t,x,z)$
- ► Form the horizontal Darcy law:  $u = -\overline{\kappa}\nabla(g\eta + \mu q)$ we get  $h\overline{u} = -g\overline{\kappa}h\nabla(h+S) - \mu\overline{\kappa}(\nabla(h\overline{q}) + \beta q_S\nabla S) + O(\mu^2)$ with  $\overline{q}(t,x) = \frac{1}{h(t,x)}\int_{S}^{\eta(t,x)}q(t,x,z) dz$  and  $q_S(t,x) = q(t,x,S(x))$ .
- ► Form the incompressibility constraint: we get  $w(t,x,z) = \overline{w}(t,x) + \frac{z-S-\frac{h}{2}}{h} \widetilde{w}(t,x) + O(\mu^2)$ with  $\overline{w} = \mu \left(\beta \overline{u} \cdot \nabla S - \frac{h}{2\overline{\Theta}} \nabla \cdot \left(\overline{\Theta} \overline{u}\right)\right)$ and  $\widetilde{w} = -\frac{h}{\overline{\Theta}} \nabla \cdot \left(\overline{\Theta} \overline{u}\right)$
- ► Form the vertical Darcy law: we get  $h\overline{w} = \overline{\kappa}q_S$  and  $h\widetilde{w} = 6\overline{\kappa}(2\overline{q} - q_S)$
- ► We conclude the equation on the horizontal velocity:  $(I_{d} + \mu^{2} \mathcal{T}_{S,h}^{\overline{\kappa},\overline{\Theta}}) \overline{u} = -g\overline{\kappa}\nabla(h+S)$ with  $\mathcal{T}_{\mathbf{b},\mathbf{h}}^{\overline{\kappa},\overline{\Theta}}(V) = \beta^{2}\nabla\mathbf{b}\otimes\nabla\mathbf{b}V + \frac{\overline{\kappa}}{\mathbf{h}}\nabla\left(\frac{\beta\mathbf{h}^{2}\nabla\mathbf{b}}{2\overline{\kappa}}\cdot V\right) \frac{\beta\mathbf{h}^{2}\nabla\mathbf{b}}{2\overline{\Theta}\mathbf{h}}\nabla\cdot\left(\overline{\Theta}V\right) + \frac{\overline{\kappa}}{\mathbf{h}}\nabla\left(\frac{\mathbf{h}^{3}}{3\overline{\kappa}\overline{\Theta}}\nabla\cdot\left(\overline{\Theta}V\right)\right).$ [Lannes'13]

<u>HYDRODYNAMIC DUPUIT-FORCHHEIMER MODEL (HDF<sub>S,h</sub>):</u>  $O(\mu^2)$ -approximate model  $\partial_t (\overline{\Theta}h) - \nabla \cdot \left(\overline{\Theta}h \left(I_d + \mu^2 \mathcal{T}_{S,h}^{\overline{K},\overline{\Theta}}\right)^{-1} (g\overline{\kappa}\nabla(h+S))\right) = 0$  [Bucur'16]

▶ Energy dissipation:  $\partial_t \mathscr{P}(h) + \nabla \cdot \mathscr{G}(h, u, w) = -\widetilde{\mathscr{K}}(h, u, w)$ 

where 
$$\mathscr{P}(h) := g\overline{\Theta}h\left(S + \frac{h}{2}\right)$$
  
 $\mathscr{G}(h, u, w) := (g(h+S) + \mu q)\overline{\Theta}h\overline{u}$   
and  $\widetilde{\mathscr{K}}(h, u, w) := \frac{\overline{\Theta}h}{\overline{\kappa}}\left(|\overline{u}|^2 + \overline{w}^2 + \frac{\widetilde{w}^2}{12}\right)$ 

► Linear decay rate: 
$$\tilde{\lambda} = \frac{|kD|^2}{1 + \frac{|kD|^2}{3}} \xrightarrow{|kD| \to \infty} 3.$$









Numerical scheme:

$$\begin{split} h_{k}^{n+1} &= h_{k}^{n} - \frac{\delta_{t}}{\overline{\Theta}_{k}} \nabla_{k}^{\delta} \cdot \left(\overline{\Theta}_{\star} \left\{h_{\star}^{n}\right\}_{\star} \overline{u}_{\star}^{n+1} \\ 1 + \left[\mathcal{T}_{S_{\star},h_{\star}^{n}}^{\overline{\kappa},\overline{\Theta}}\right]_{f} \right) \left(\overline{u}_{\star}^{n+1}\right) &= -g\overline{\kappa}_{f} \nabla_{f}^{\delta} \left(h_{\star}^{n+1} + S_{\star}\right). \end{split}$$

The discrete space differential operator is defined as

$$\begin{split} \left[ \mathcal{T}_{\mathbf{b}_{\star},\mathbf{h}_{\star}}^{\overline{\kappa},\overline{\Theta}} \right]_{f} \left( U_{\star} \right) &= \frac{\overline{\kappa}_{f}}{\{\mathbf{h}_{\star}\}_{f}} \left\{ \frac{\mathbf{h}_{\star}}{\overline{\kappa}_{\star}\overline{\Theta}_{\star}} \left\{ \overline{\Theta}_{\star} U_{\star} \cdot \nabla_{\star}^{\delta} \mathbf{b}_{\star} \right\}_{\star} \right\}_{f} \nabla_{f}^{\delta} \mathbf{b}_{\star} \\ &+ \frac{\overline{\kappa}_{f}}{\{\mathbf{h}_{\star}\}_{f}} \left( \nabla_{f}^{\delta} \left( \frac{|\mathbf{h}_{\star}|^{2}}{2\overline{\kappa}_{\star}\overline{\Theta}_{\star}} \left\{ \overline{\Theta}_{\star} U_{\star} \cdot \nabla_{\star}^{\delta} \mathbf{b}_{\star} \right\}_{\star} \right) - \left\{ \frac{|\mathbf{h}_{\star}|^{2}}{2\overline{\kappa}_{\star}\overline{\Theta}_{\star}} \nabla_{\star}^{\delta} \cdot \left( \overline{\Theta}_{\star} U_{\star} \right) \right\}_{f} \nabla_{f}^{\delta} \mathbf{b}_{\star} \right) \\ &- \frac{\overline{\kappa}_{f}}{\{\mathbf{h}_{\star}\}_{f}} \nabla_{f}^{\delta} \left( \frac{|\mathbf{h}_{\star}|^{3}}{3\overline{\kappa}_{\star}\overline{\Theta}_{\star}} \nabla_{\star}^{\delta} \cdot \left( \overline{\Theta}_{\star} U_{\star} \right) \right). \end{split}$$

The discrete divergence and gradient operators are defined as

$$\begin{aligned} \nabla_k^{\delta} \cdot U_{\star} &= \frac{1}{\mathfrak{m}_k} \sum_{f \in \mathbb{F}_k} U_f \cdot \mathfrak{n}_k^{k_f} \mathfrak{m}_f \quad \text{and} \quad \nabla_f^{\delta} \phi_{\star} &= \frac{\phi_{k_f} - \phi_k}{\delta_f} \mathfrak{n}_k^{k_f} \\ & \{\psi_{\star}\}_k &= \frac{1}{\#\mathbb{F}_k} \sum_{f \in \mathbb{F}_k} \psi_f \quad \text{and} \quad \{\phi_{\star}\}_f &= \frac{\phi_{k_f} + \phi_k}{2}. \end{aligned}$$

Entropy satisfying: The scheme satisfies a discrete counterpart of the energy dissipation.











## Water Waves model with drag force $(WW_{\kappa,\Theta})$ :



#### PROPERTIES

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► Energy dissipation: 
$$\partial_t \left( \mathscr{P}(\eta) + \int_S^{\eta} \Theta \mathscr{K}(u, w) \, dz \right) + \nabla \cdot \mathscr{G}(\eta, u, w) = -\widetilde{\mathscr{K}}(\eta, u, w)$$
  
where  $\mathscr{P}(\eta) := g \int_S^{\eta} \Theta z \, dz$   
 $\mathscr{K}(u, w) := \frac{1}{2} \left( |u|^2 + w^2 \right)$   
 $\mathscr{G}(\eta, u, w) := \int_S^{\eta} \left( gz + \mathscr{K}(u, w) + p \right) \Theta u \, dz$   
and  $\widetilde{\mathscr{K}}(\eta, u, w) := \int_S^{\eta} \frac{\Theta}{\kappa} \left( |u|^2 + w^2 \right) \, dz.$ 

 $\kappa \ll 1, \ \Theta \approx 10^{-1}$ 

Assuming hydrostatic pressure and vertical-homogeneous horizontal velocity.

▶ By vertical integration on  $[S, \min(B, \eta)]$  and  $[\min(B, \eta), \eta]$ .



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- > Assuming hydrostatic pressure and vertical-homogeneous horizontal velocity.
- ▶ By vertical integration on  $[S, \min(B, \eta)]$  and  $[\min(B, \eta), \eta]$ .

#### PROPERTIES

 $\begin{array}{l} & \overbrace{\partial_{t}\mathcal{V}(\eta) + \nabla \cdot (\mathcal{H}^{s}(\eta)\overline{u}^{s})}_{\partial_{t}\mathcal{V}(\eta) + \nabla \cdot (\mathcal{H}^{s}(\eta)\overline{u}^{s})} = \nabla \cdot (g\mathcal{V}^{g}(\eta)\overline{\kappa}^{g}(\eta)\nabla\eta) \\ & \overbrace{\partial_{t}(\mathcal{H}^{s}(\eta)\overline{u}^{s}) + \nabla \cdot (\mathcal{H}^{s}(\eta)\overline{u}^{s}\otimes\overline{u}^{s})}_{\partial t} = -g\mathcal{H}^{s}(\eta)\nabla\eta + \frac{\partial_{t}\mathcal{V}^{s}(\eta) + \nabla \cdot (\mathcal{V}^{s}(\eta)\overline{u}^{s})}{2}\overline{u}^{s} \\ & \text{with} \quad \mathcal{H}^{s}(\eta) := \max(0, \eta - B) \\ \end{array}$ 

► Energy dissipation:  $\partial_t (\mathscr{P}(\eta) + \mathscr{K}(\eta, \overline{u}^s)) + \nabla \cdot \mathscr{G}(\eta, \overline{u}^s) = -\widetilde{\mathscr{K}}(\eta)$ where  $\mathscr{P}(\eta) := g \int_S^{\eta} \Theta z \, dz$ ,  $\mathscr{K}(\eta, \overline{u}^s) := \frac{\mathscr{K}^s(\eta)}{2} |\overline{u}^s|^2$   $\mathscr{G}(\eta, \overline{u}^s) := (g\eta \mathscr{H}^s(\eta) + \mathscr{K}(\eta, \overline{u}^s)) \overline{u}^s - g^2 \eta \mathscr{V}^g(\eta) \overline{\kappa}^g \nabla \eta$ and  $\widetilde{\mathscr{K}}(\eta) := g^2 \mathscr{V}^g(\eta) \overline{\kappa}^g |\nabla \eta|^2$ .



$$\frac{\partial_{t} \mathcal{V}(\eta) + \nabla \cdot (\mathcal{H}^{s}(\eta) \overline{u}^{s})}{\partial_{t} (\mathcal{H}^{s}(\eta) \overline{u}^{s}) + \nabla \cdot (\mathcal{H}^{s}(\eta) \overline{u}^{s}) = \nabla \cdot (g \mathcal{V}^{g}(\eta) \overline{\kappa}^{g}(\eta) \nabla \eta)}{\partial_{t} (\mathcal{H}^{s}(\eta) \overline{u}^{s}) + \nabla \cdot (\mathcal{H}^{s}(\eta) \overline{u}^{s} \otimes \overline{u}^{s}) = -g \mathcal{H}^{s}(\eta) \nabla \eta + \frac{\partial_{t} \mathcal{V}^{s}(\eta) + \nabla \cdot (\mathcal{V}^{s}(\eta) \overline{u}^{s})}{2} \overline{u}^{s}}$$
with  $\mathcal{H}^{s}(\eta) := \max(0, \eta - B)$ 
ROPERTIES

► Energy dissipation: 
$$\partial_t (\mathscr{P}(\eta) + \mathscr{K}(\eta, \overline{u}^s)) + \nabla \cdot \mathscr{G}(\eta, \overline{u}^s) = -\widetilde{\mathscr{K}}(\eta)$$
  
where  $\mathscr{P}(\eta) := g \int_S^{\eta} \Theta z \, dz$ ,  $\mathscr{K}(\eta, \overline{u}^s) := \frac{\mathscr{H}^s(\eta)}{2} |\overline{u}^s|^2$   
 $\mathscr{G}(\eta, \overline{u}^s) := (g\eta \mathscr{H}^s(\eta) + \mathscr{K}(\eta, \overline{u}^s)) \overline{u}^s - g^2 \eta \mathscr{V}^g(\eta) \overline{\kappa}^g \nabla \eta$   
and  $\widetilde{\mathscr{K}}(\eta) := g^2 \mathscr{V}^g(\eta) \overline{\kappa}^g |\nabla \eta|^2$ .

$$\underbrace{ \text{Unified Green-Naghdi/Hydrodynamic Dupuit-Forchheimer model } (GN/HDF):}_{\partial_t \mathcal{V}(\eta) + \nabla \cdot (\mathcal{H}^s(\eta)\overline{u}^s) = \nabla \cdot \left(g\mathcal{V}^g(\eta) \left(I_d + \mathcal{T}_{S,\mathcal{H}^g(\eta)}^{\overline{\kappa},\overline{\Theta}}\right)^{-1} (\overline{\kappa}^g(\eta)\nabla\eta)\right)}_{\left(I_d + \mathcal{T}_{B,\mathcal{H}^s(\eta)}^{1,1}\right)^{-1} \partial_t \overline{u}^s + \overline{u}^s \cdot \nabla \overline{u}^s = -g\mathcal{H}^s(\eta)\nabla\eta - Q_{B,\eta}(\overline{u}^s)}$$



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Numerical scheme: ImEx scheme

[Parisot,Vila'16] 👔 [Lteif,Kazolea,Parisot'25]

Implicit continuity equation

 $\mathcal{V}\left(\eta_{k}^{n+1}\right) = \mathcal{V}\left(\eta_{k}^{n}\right) - \delta_{t} \nabla_{k}^{\delta} \cdot \left(\left\{\mathcal{H}^{s}\left(\eta_{\star}^{n+1}\right) \overline{u}_{\star}^{n}\right\}_{\star} - \left(\delta_{t}\left\{\gamma\left(\eta_{\star}^{n+1}\right)\right\}_{\star} + \left\{g\mathcal{V}^{g}\left(\eta_{\star}^{n+1}\right) \overline{\kappa}_{\star}^{g}\right\}_{\star}\right) \nabla_{\star}^{\delta}\left(\eta_{\star}^{n+1}\right)\right)$ 

Explicit vertical exchanges

$$G_{k}^{n+1} = \frac{H_{k}^{s*} - \mathcal{H}^{s}\left(\eta_{k}^{n+1}\right)}{\delta_{t}} \quad \text{with} \quad H_{k}^{s*} = \mathcal{H}^{s}\left(\eta_{k}^{n}\right) - \delta_{t}\nabla_{k}^{\delta} \cdot \left(F_{\star}^{n+1}\right)$$
  
and 
$$F_{f}^{n+1} = \left\{\mathcal{H}^{s}\left(\eta_{\star}^{n+1}\right)\overline{u}_{\star}^{n}\right\}_{f} - \delta_{t}\left\{\gamma\left(\eta_{\star}^{n+1}\right)\right\}_{f}\nabla_{f}^{\delta}\left(\eta_{\star}^{n+1}\right)$$

► Explicit momentum equation (si  $\mathcal{V}^{s}(\eta_{k}^{n+1}) > 0$ )

$$\left(\mathcal{V}^{s}\left(\eta_{k}^{n+1}\right)+\frac{G_{k}^{n+1}}{2}\right)\overline{u}_{k}^{n+1}=\mathcal{V}^{s}\left(\eta_{k}^{n}\right)\overline{u}_{k}^{n}-\delta_{t}\left(\nabla_{k}^{\mathrm{up}}\cdot\left(\overline{u}_{\star}^{n},F_{\star}^{n+1}\right)+g\mathcal{H}^{s}\left(\eta_{k}^{n+1}\right)\nabla_{k}^{\delta}\left(\eta_{\star}^{n+1}\right)\right)$$

**Entropy satisfying:** for  $\gamma(\eta) \ge g\mathcal{V}^{s}(\eta)$  and under a CFL condition, the scheme satisfies a discrete counterpart of the energy dissipation.









### Conclusion:

- Derivation of an approximate hydrodynamic model improving the linear decay rate preserving the energy dissipation
- Unified model for shallow water and Dupuit-Forchheimer regime "coupling" between hyperbolic and elliptic models preserving the energy dissipation



[Carreau,Parisot'25]

## Perspectives:

- ▶ Analysis of the hydrodynamic model (maximum principle, monotony, TVD...)
- ▶ Analysis of the numerical stability (with/without CFL condition)
- Higher order hydrodynamic models
- ▶ Vertical dependency of the physical parameters  $\kappa$  and  $\Theta$
- Infiltration model in the vadose area
- ► Study of interface conditions T. Bonnet, M. Coquerelle (I2M)
- Morphodynamic coupling (physical parameters evolution, deformation)
- Same questions for the Richards model (scalar 3D advection-diffusion equation)

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## Thank You